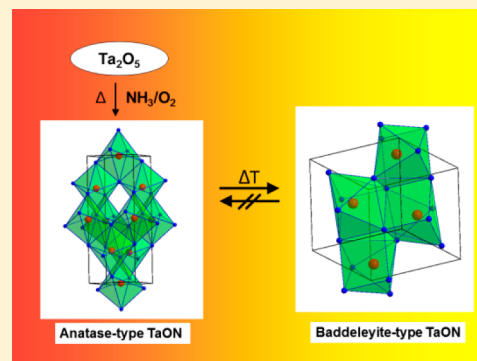


Synthesis and Crystal Structure of δ -TaON, a Metastable Polymorph of Tantalum Oxide NitrideTobias Lüdtkke,[†] Alexander Schmidt,[†] Caren Göbel,[†] Anna Fischer,[†] Nils Becker,^{‡,§} Christoph Reimann,^{||} Thomas Bredow,^{||} Richard Dronskowski,^{‡,§} and Martin Lerch^{*,†}[†]Institut für Chemie, TU Berlin, Straße des 17. Juni 135, D-10623 Berlin, Germany[‡]Institut für Anorganische Chemie, RWTH Aachen University, Landoltweg 1, D-52056 Aachen, Germany[§]Jülich-Aachen Research Alliance (JARA-HPC), RWTH Aachen University, D-52056 Aachen, Germany^{||}Institut für Physikalische und Theoretische Chemie, Universität Bonn, Wegelerstraße 12, D-53115 Bonn, Germany

Supporting Information

ABSTRACT: δ -TaON was prepared by reaction of gaseous ammonia with an amorphous tantalum oxide precursor. As a representative of the anatase structure (aristotype) it crystallizes in the tetragonal crystal system with lattice parameters $a = 391.954(16)$ pm and $c = 1011.32(5)$ pm. At temperatures between 800 and 850 °C an irreversible phase transformation to baddeleyite-type β -TaON is observed. While quantum-chemical calculations confirm the metastable character of δ -TaON, its transformation to β -TaON is kinetically controlled. The anion distribution of the anatase-type phase was studied theoretically. In agreement with previous studies, it was found that a configuration with maximal N–N distances is most stable. The calculated band edge energies indicate that δ -TaON is a promising photocatalytic material for redox reactions, e.g., water splitting.



INTRODUCTION

Transition metal oxide nitrides are a group of materials which attracted a lot of attention in recent years. In particular, tantalum-based compounds are, for example, suitable as nontoxic color pigments,^{1,2} can be considered as fast ion conductors,³ and are used as materials for photocatalytic water splitting under sunlight.^{4,5} From a thermodynamical point of view, the most stable polymorph of tantalum oxide nitride is the so-called β -phase, reported first by Brauer and Weidlein.⁶ It crystallizes in the well-known monoclinic baddeleyite-type structure also reported for ZrO_2 and exhibits a bright yellow color (optical band gap of ~ 2.6 eV). A metastable polymorph (γ -TaON) was reported by Schilling et al.⁷ It shows an orange color (optical band gap of ~ 2.15 eV) and crystallizes similar to $VO_2(B)$ in the monoclinic crystal system. For both β - and γ -TaON, an ordered distribution of oxygen and nitrogen was found. Another polymorph, α -TaON, was suggested by Buslaev et al.⁸ but disproved by quantum-chemical calculations.⁹ A possible high-pressure modification was predicted independently by two different groups.^{10,11} Pressures of more than ~ 31 GPa should induce a phase transformation from β -TaON (7-fold coordination of Ta) to a cotunnite-type structure with 9-fold coordinated cations. In fact, this cotunnite phase has recently been found by high-pressure studies,¹² just as predicted. In addition, the relative stabilities of various structure types (including the anatase, rutile, and fluorite type) were calculated by quantum-chemical methods a couple of years ago.¹³ Interestingly, and in clear

contrast to TiO_2 , the anatase type is energetically favored compared to the rutile type. With the small calculated stability difference of only ~ 20 kJ/mol between baddeleyite (most stable) and anatase type,¹³ it is worth the effort trying to synthesize anatase-type TaON.

It should be mentioned that anatase-type phases have previously been prepared by incorporating lower-valent cations such as Mg^{2+14} or Sc^{3+15} into TaON. In particular, substitution of 5% tantalum by magnesium leads to a phase-pure material with the composition $Mg_{0.05}Ta_{0.95}O_{1.15}N_{0.85}$. This was the first example of an anatase-type structure not containing the element titanium. Interestingly, the preparation of 30–45 nm thin films of undoped anatase-type TaON on $(LaAlO_3)_{0.37}(SrAl_{0.5}Ta_{0.5}O_3)_{0.7}$ (LSAT) single crystals using nitrogen plasma-assisted pulsed laser deposition (NPA-PLD) was recently reported by Suzuki et al.¹⁶ Epitaxial stress from the substrate stabilizes the anatase-type structure without additional elements such as magnesium or scandium. From a photocatalytic point of view, this material is of great interest due to its optical band gap of ~ 2.4 eV and its high charge carrier mobility which is comparable to that of anatase-type TiO_2 .¹⁶ Unfortunately, the quality of the X-ray diffraction patterns of the films is insufficient for a real structural refinement including atomic positions, bond lengths, etc. In the present Article we report, for the first time, the synthesis and crystal structure of

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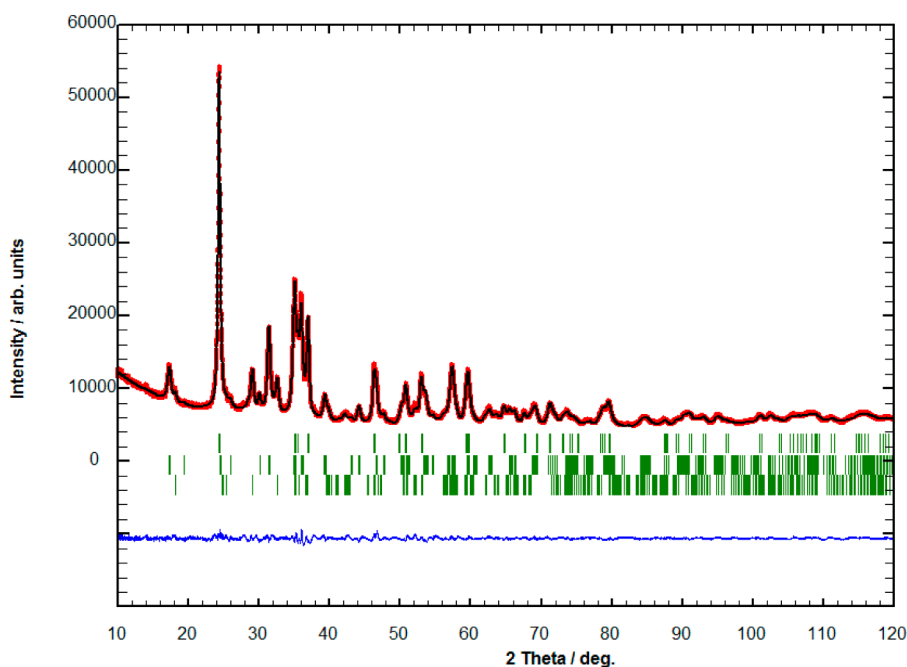


Figure 1. X-ray powder diffraction pattern with results of the Rietveld refinement. The vertical bars indicate the reflection positions of anatase-type δ -TaON (top), Ta_3N_5 (middle), and β -TaON (bottom).

tantalum oxide nitride powder with anatase-type structure without any additional cations. From now on we call this the δ -TaON phase.

EXPERIMENTAL SECTION

Synthesis. Amorphous tantalum oxide precursors were prepared via a modified Pechini route.¹⁷ Tantalum chloride (99.999%, Alfa Aesar) was solved in an ethanol–citric acid solution. Citric acid was used in a molar excess of 12 times the TaCl_5 . Ethylene glycol in a molar excess of 17 times the tantalum chloride was added for the gelling process. Ethanol, HCl, and water, the latter formed during the esterification, are evaporated at $\sim 150^\circ\text{C}$. The polymerization can be finalized by heating the mixture to 200°C . The material was calcined at 400°C for 40 h. The products were colorless X-ray amorphous powders. The maximum yield of 43 wt % of anatase-type tantalum oxide nitride was obtained by ammonolysis of these amorphous precursors at 760°C , a reaction time of 1 h, and flow rates of 12.50 L h^{-1} ammonia and 0.02 L h^{-1} oxygen. Additional phases were β -TaON, γ -TaON, and Ta_3N_5 . Samples without γ -phase could be obtained using an increased ammonia flow rate of 23.75 L h^{-1} , while the other parameters were unaltered. For the investigation of the high-temperature behavior in inert atmosphere the samples were heated at different temperatures for 1 h in closed copper capsules (nitrogen atmosphere). After cooling down to ambient temperature the phase fractions were determined by X-ray diffraction.

Characterization. A PANalytical X'Pert Pro powder diffractometer (Cu $K\alpha$ radiation) was used for XRD measurements. Rietveld refinements including quantitative phase analyses were carried out using the program FULLPROF 2000.¹⁸ Profiles were fitted with a pseudo-Voigt function. The contents of nitrogen and oxygen were determined by hot gas extraction using a LECO TC-300/EF-300 N/O analyzer. ZrO_2 and steel were used as standard materials. The accuracy is $\sim 2\%$ of the N/O contents present. X-ray fluorescence spectrometry (XRF) was performed using a PANalytical Axios spectrometer with an Rh-tube. Crystallinity, phase composition (identification of δ -TaON, β -TaON, Ta_3N_5), and morphology of the particles were investigated by high-resolution transmission electron microscopy (HRTEM), selected area electron diffraction (SAED), and energy dispersive X-ray spectroscopy (EDX) analysis. For TEM analysis, powder samples were deposited on a holey carbon coated nickel grid. TEM and

HRTEM measurements were performed at the Department of Electron Microscopy of the TU Berlin (ZELMI) on a Tecnai G² S-Twin transmission electron microscope (FEI Company, Eindhoven, Netherlands) equipped with a LaB₆-source operated at 200 kV. Images were recorded with a GATAN MS794 P CCD-camera. EDX-analysis were carried out with an r-TEM SUTW Detector from EDAX (Si (Li)-detector).

Quantum-Chemical Calculations. In order to investigate the thermodynamic properties of the TaON phases, periodic density-functional theory (DFT) calculations were performed as implemented in the Vienna ab initio simulation package (VASP).^{19,20} Projector-augmented waves (PAW)²¹ were used to separate the core and valence electrons. Exchange and correlation contributions were treated within the generalized-gradient-approximation as described by Perdew, Burke, and Ernzerhof (GGA-PBE).²² An energy-cutoff of 500 eV and a Monkhorst–Pack k -point mesh of $20 \times 20 \times 8$ were used for structure optimization, and ensured convergence.

Quasiharmonic phonon calculations were performed using the program Phonopy.²³ In a first step, a supercell was constructed on the basis of the optimized crystal structure. All symmetry-inequivalent atoms of the previously created supercell were then slightly shifted out of their equilibrium position to calculate the Hellman–Feynman forces. Subsequently, the force constants and then the dynamical matrix were computed. With a combination of its eigenvalues (phonon frequencies) with Bose–Einstein statistics, the free phonon energy was obtained.²⁴ Thermodynamic properties at finite temperatures are straightforwardly accessible from these data.

For the investigation of the high-pressure behavior, the cell volume was varied by scaling the cell parameters between 0.94 and 1.04 (with a fixed $a:b:c$ ratio and rectangular cell angles). The resulting energy versus volume data were fitted to the Birch–Murnaghan equation of state.²⁵ Thus, the pressure and, eventually, reaction enthalpies as a function of the pressure were calculated.

In order to investigate the dependence of the calculated properties on the basis set and functional, additional periodic calculations employing the PBE functional as well as the PBE0²⁶ DFT–Hartree–Fock hybrid method were conducted with the CRYSTAL program package (C09).²⁷ We used atom-centered Gaussian basis sets of valence double- ζ quality that have been optimized especially for the use in solid-state calculations (exponents and coefficients are listed in the Supporting Information). For tantalum an effective core potential

Table 1. Results of the Rietveld Refinement for δ -TaON in Comparison with Selected Results for $\text{Mg}_{0.05}\text{Ta}_{0.95}\text{O}_{1.15}\text{N}_{0.85}$

	δ -TaON	$\text{Mg}_{0.05}\text{Ta}_{0.95}\text{O}_{1.15}\text{N}_{0.85}$ ¹³
structure type	anatase	anatase
space group	$I4_1/amd$	$I4_1/amd$
cryst syst	tetragonal	tetragonal
lattice params	$a = 391.954(16)$ pm $c = 1011.32(5)$ pm $c/a = 2.58$	$a = 391.986(6)$ pm $c = 1011.19(3)$ pm $c/a = 2.58$
unit cell volume	$V = 155.367(12) \times 10^6$ pm ³	$V = 155.373(6) \times 10^6$ pm ³
formula units	$Z = 4$	$Z = 4$
calcd density	$\rho = 9.02$ g/cm ³	$\rho = 8.71$ g/cm ³
diffractometer	PANalytical X'Pert Pro	
2θ range	10–120°	
wavelength	Cu $K\alpha 1$, $\lambda = 154.053$ pm Cu $K\alpha 2$, $\lambda = 154.439$ pm	
profile points	8839	
R_{wp}	2.02%	
R_{exp}	1.15%	
R_{Bragg}	1.41%	
S	1.75	

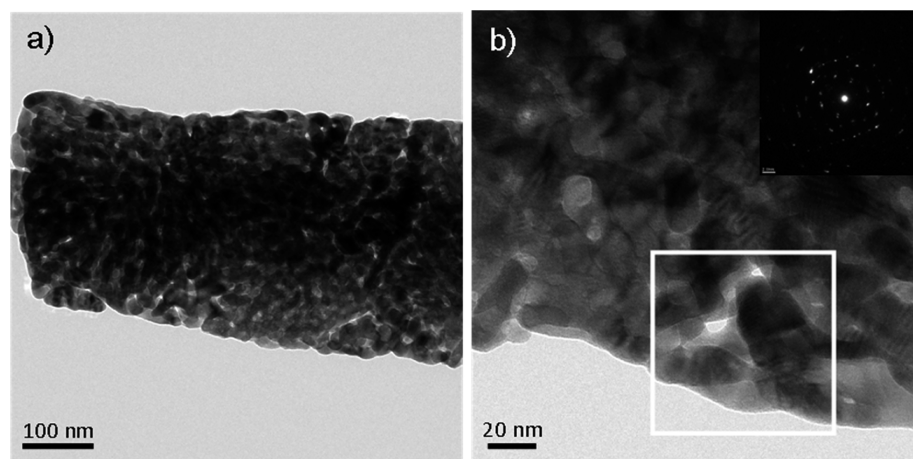


Figure 2. TEM images of the as-synthesized sample containing δ -TaON, β -TaON, and Ta_3N_5 : (a, b) porous sponge-like structure. SAED pattern given as inset, demonstrating the presence of multiple crystal phases.

was employed (SD60MWB²⁸) so that a total number of 13 electrons were kept in the semicore and valence region of each Ta atom, while all-electron basis sets were used for oxygen and nitrogen. The default values for the integral accuracy in CRYSTAL09 have been increased by a factor of 10, and Monkhorst–Pack and Gilat grids have been generated using the shrinking factors (8,8). In order to obtain the Gibbs energies of the TaON polymorphs we performed frequency calculations at the Γ point.^{29,30} Furthermore, we calculated the bulk modulus for the new phase by fitting the results of constant volume optimizations against the third-order isothermal Birch–Murnaghan’s equation of state.

For the determination of the absolute electronic band edge positions with respect to the vacuum we performed periodic slab calculations with the PBE as well as the PBE0 functional. Since the Kohn–Sham gap as obtained from the PBE calculations is not a good approximation to the optical gap in solids,³¹ we will follow the approach of Caspary Toroker et al. and merely calculate the band gap center from the hybrid calculations.³² It has been shown that the band gap center is reproduced more reliably within the Kohn–Sham approach than the valence band maximum.³³

In order to calculate the optical excitations, we employed the Bethe–Salpeter approach.³⁴ Therefore, we first calculated the quasiparticle gap with the GW0 method using the VASP code.^{35–38} The computational setup for these surface calculations has been described

in detail elsewhere.³⁹ For these calculations, we used small-core PAWs developed for GW calculations as supplied with the VASP program package. In order to reduce the computational demands, a smaller Monkhorst–Pack k -point grid ($6 \times 6 \times 6$) has been chosen as well as a lower energy cutoff (400 eV) compared to the structure relaxations. Increasing these values in test calculations did not alter the results significantly.

RESULTS AND DISCUSSION

The synthetic conditions, described in the Experimental Section, had to be respected carefully. Slightly increased reaction temperatures (e.g., 770 °C instead of 760 °C) strongly preferred the formation of Ta_3N_5 , while at lower temperatures (for example 750 °C) large parts of the amorphous tantalum oxide precursor did not react with ammonia but partially crystallized to β - Ta_2O_5 . As mentioned above, we were successful in obtaining samples with a yield up to 43 wt % of the new δ -TaON polymorph as best result. With respect to the difficulties in refining the crystal structure of the disordered γ -phase,⁷ we decided to choose a sample with a slightly smaller fraction of δ -TaON but without any γ -TaON for the Rietveld refinement. The red color of all prepared samples is likely

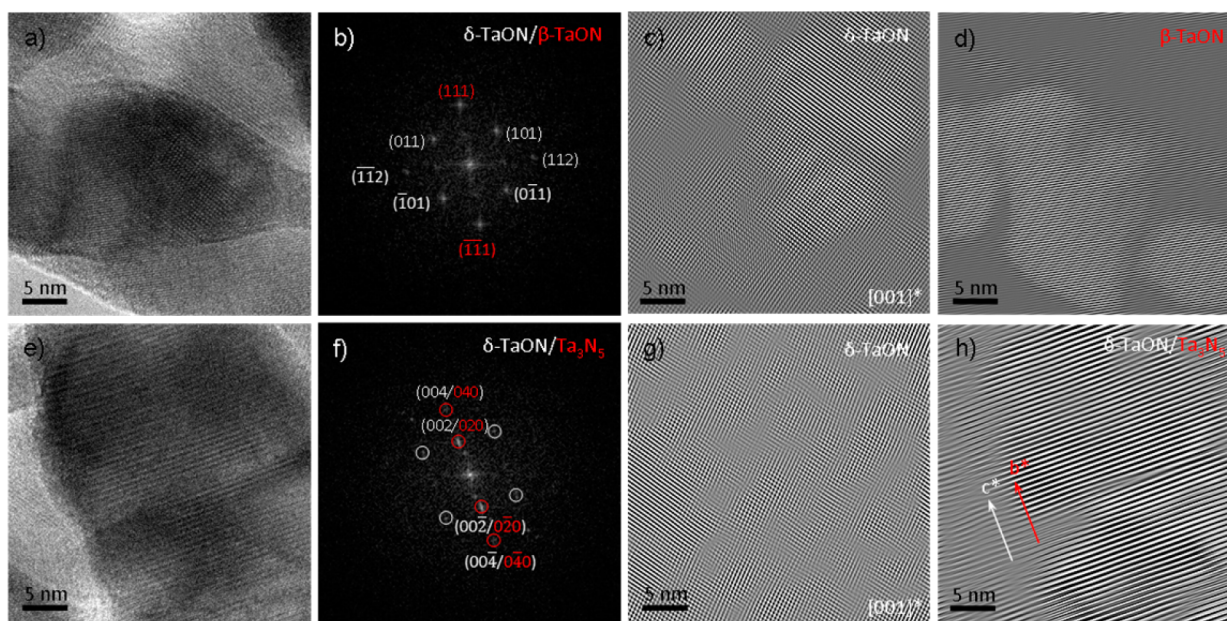


Figure 3. HRTEM-images (a, e) from the marked region in Figure 2 with corresponding FFT patterns (b, f). Filtered HRTEM-images of (a) evidencing the presence of δ -TaON (c) and β -TaON (d) and filtered HRTEM-images (e) evidencing the presence of δ -TaON (g) and Ta_3N_5 (h).

dominated by the presence of dark-red Ta_3N_5 . The red color also points to the absence of significant amounts of reduced tantalum resulting in black samples. XRF analysis proved the absence of other elements than tantalum, oxygen, and nitrogen.

Figure 1 presents the experimental powder XRD pattern with the results of the Rietveld refinement. Details are given in Table 1. The sample consists of three crystalline phases: anatase-type δ -TaON (35 wt %), Ta_3N_5 (45 wt %), and baddeleyite-type β -TaON (20 wt %). This is in good agreement with the results of chemical analysis (N = 7.6 wt %, O = 6.1 wt %), taking into account an additional content of amorphous tantalum oxide (assuming an overall content of 13 wt % amorphous Ta_2O_5 for the calculation gives the values from chemical analysis).

Additional TEM-investigations revealed that the sample consists of agglomerated crystalline nanoparticles assembled in sponge-like structures (Figure 2). Electron diffraction at those agglomerates as well as HRTEM analysis (see filtered images and corresponding FFT) of single nanoparticles revealed the presence of homogeneously intergrown δ -TaON, Ta_3N_5 , and β -TaON particles on the nanometer scale (20–50 nm particle size, Figure 3). While δ -TaON and Ta_3N_5 are present as rounded “sintered” particles, β -TaON tends also to form angular aggregates. Unfortunately EDX analysis of single particles of identified phases could not be performed, due to intergrowth/overlapping of the nanoparticles. The observed overall ratio N/O of 1.23 is in very good agreement with the above-presented results of the chemical analysis (N/O = 1.24).

In comparison to $\text{Mg}_{0.05}\text{Ta}_{0.95}\text{O}_{1.15}\text{N}_{0.85}$, the lattice parameters determined for δ -TaON have very similar values (see Table 1). This can be easily explained by the similar ionic radii and the small amount of Mg^{2+} in the doped phase. In contrast, the calculated density of the undoped material is significantly larger due to the low mass of magnesium. It is remarkable that the density of the more stable γ -TaON phase (8.64 g/cm³) is clearly smaller which comes from the special structural features of the $\text{VO}_2(\text{B})$ type (huge tunnels, etc., see ref 7).

The refined atomic parameters are presented in Table 2. The isotropic Debye–Waller factors were fixed to the results

Table 2. Refined Atomic Parameters for Anatase-type δ -TaON

atom	Wyckoff	x	y	z	SOF ^a	$B_{\text{iso}}/\text{\AA}^2$
Ta	4b	0	1/4	3/8	1	1.3
O/N ^b	8e	0	1/4	0.5789(7)	0.5/0.5	1.4

^aSite occupation factor. ^bNot refined independently.

acquired from neutron diffraction experiments on the Mg-containing phase.¹⁴ Anatase-type TiO_2 , which is structurally characterized by TiO_6 -octahedra sharing four edges with neighboring octahedra and channels along the *a*- and *b*-axes, crystallizes in the tetragonal crystal system in space group $I4_1/amd$ (No. 141). Using this space group for the refinement of δ -TaON does not give the opportunity to respect an ordered arrangement of nitrogen and oxygen. In our case this is not a real problem because X-ray diffraction experiments do not allow distinguishing between these elements (ions with the same number of electrons). Figure 4 depicts the crystal structure of anatase-type TaON with the above-mentioned aristotype setting with just one kind of anion position. The calculated bond lengths of the octahedra are also presented in Figure 4 and are in good agreement with the ionic radii expected from the corresponding ions⁴⁰ and that of $\text{Mg}_{0.05}\text{Ta}_{0.95}\text{O}_{1.15}\text{N}_{0.85}$.¹⁴

Possible anion arrangements have already been studied theoretically on a quantum-chemical level¹³ and experimentally using $\text{Mg}_{0.05}\text{Ta}_{0.95}\text{O}_{1.15}\text{N}_{0.85}$ ^{14,41} with consistent results. There are three maximal nonisomorphic subgroups (all *translationengleich*) providing two nonequivalent anion sites: $I\bar{4}m2$ (No. 119), $I4_1md$ (No. 109), and $Imma$ (No. 74). It was shown that space group $I4_1md$ is the most preferred for tantalum oxide nitrides with anatase-type structure. Unfortunately, the amount of material (~100 mg) is hitherto not sufficient for a neutron diffraction investigation which is the method of choice for the experimental determination of the N/O distribution. Larger amounts of material and an increase of the fraction of the δ -phase are in the focus of our work in the near future.

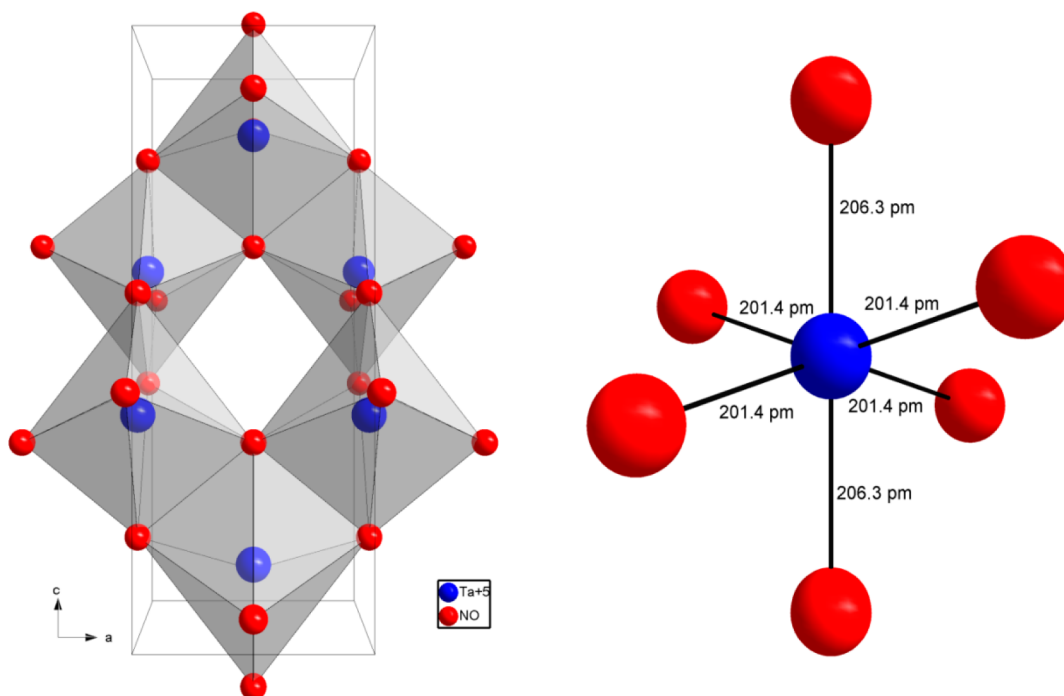


Figure 4. Crystal structure of anatase-type TaON. Left: unit cell with coordination octahedra. Right: Ta–O/N bond lengths.

The effect of different N/O configurations in the conventional unit cell of δ -TaON has already been investigated using periodic quantum-chemical calculations.¹³ Since we used different density functionals in this work, we nevertheless performed structure relaxations at the GGA (PBE) and DFT-HF hybrid (PBE0) level for different distributions of N and O atoms over the 8e Wyckoff position. With both functionals, almost identical results were obtained that confirm the previously found structure. In the most stable configuration, N–Ta–N and O–Ta–O chains in *a* or *b* direction alter along the *c* direction (see Figure 5). This crystal structure can be

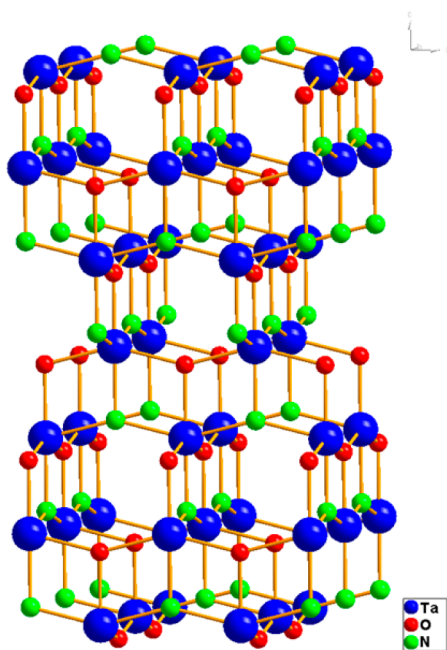


Figure 5. Supercell of δ -TaON with the most stable N/O distribution.

described in the space group $I4_1md$ (No. 109) with all atoms residing on the Wyckoff position 4a (VASP/PBE: Ta: 0,0,0.5; O: 0,0,0.299; N: 0,0,0.701).

The relative stability of the different anatase-type phases with respect to each other is linked to the N–N distance. In the most stable phase with $I4_1md$ symmetry, two N^{3-} ions are separated by at least 323 pm. In all other configurations under consideration the smallest N–N distance was reduced by more than 8% compared to this value, which leads to energy differences up to 24 kJ/mol per TaON formula unit due to the larger Coulomb repulsion between the N^{3-} anions. Nevertheless a large number of anion configurations were higher in energy by only 5–10 kJ/mol and might thus be accessible at the applied conditions during the synthesis. The theoretical Ta–O/N bond lengths differ from the experimentally assigned values as the octahedral coordination is more strongly distorted due to the reduced symmetry. At the PBE0 level we found a single short Ta–N ($d = 195$ pm) and two Ta–O bond lengths ($d = 197$ pm), two N atoms ($d = 205$ pm) at intermediate distances, and one O atom with a greater displacement ($d = 232$ pm). The slightly shorter Ta–N bond lengths as compared to Ta–O nicely reflect the larger covalency in the Ta–N combination, as also found in other studies.⁴²

In order to compare δ -TaON with the well-known β - and γ -phases, the energy per formula unit was plotted against the cell volume (Figure 6, VASP/PBE). Anatase-type TaON is energetically less favored than the thermodynamically stable baddeleyite-type by about 0.18 eV whereas the $VO_2(B)$ -type is less stable by only 0.07 eV.

From the constant volume optimizations, we calculated the bulk moduli for the three phases from the Birch–Murnaghan equation of state. The value for β -TaON ($B^{PBE} = 259$ GPa, $B^{PBE0} = 286$ GPa) is in good agreement with the results from earlier theoretical work ($G^{LDA} = 280$ GPa,⁴³ $B^{PW91} = 278$ GPa⁴⁴), whereas the previously calculated bulk modulus for the more open γ -TaON ($B^{VASP-PBE} = 181$ GPa) is less well-

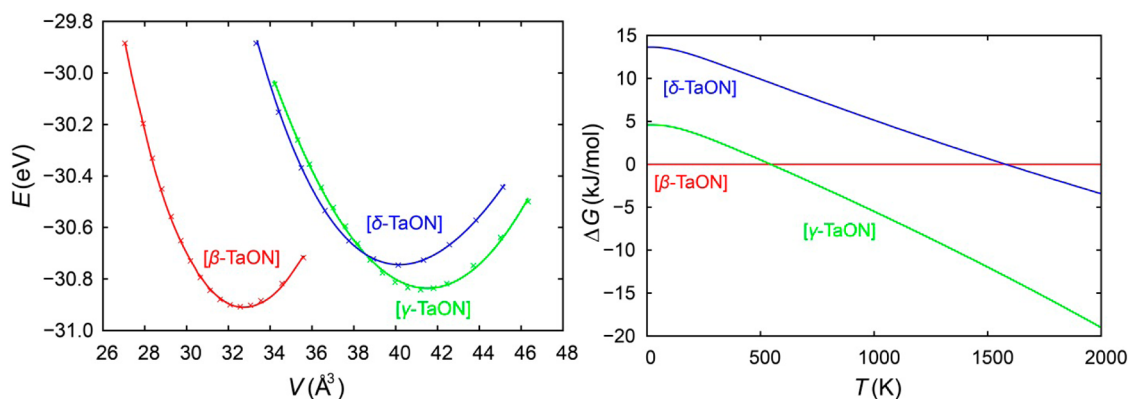


Figure 6. Energy per formula unit versus volume of the three known phases of TaON at ambient pressure (left), and relative Gibbs energy versus temperature for the different TaON phases (VASP/PBE) (right).

reproduced using atom-centered basis functions ($B^{\text{PBE}} = 210$ GPa, $B^{\text{PBE0}} = 227$ GPa). For δ -TaON we find a lower value compared to the other TaON phases ($B^{\text{PBE}} = 186$ GPa, $B^{\text{PBE0}} = 216$ GPa).

Metastable polymorphs, according to the Ostwald–Volmer rule, are usually characterized by low densities and transform to stable modifications at elevated temperatures.⁴⁵ These empirical findings are apparent in Figure 6 where δ -TaON phase has a larger volume per formula unit than β -TaON. At the same time, γ -TaON which is more stable than the δ -phase has a slightly smaller density [C09/PBE0 : $\rho(\gamma\text{-TaON}) = 8.6$ g/cm³, $\rho(\delta\text{-TaON}) = 8.9$ g/cm³]. Both values are in good agreement with the experimental results [$\rho(\gamma\text{-TaON}) = 8.6$ g/cm³, $\rho(\delta\text{-TaON}) = 9.0$ g/cm³].

In air, all nitrogen-containing phases reacted to tantalum oxide at 500 °C. In contrast, in nitrogen atmosphere anatase-type TaON undergoes a phase transition between 800 and 850 °C to the baddeleyite-type β -TaON with a density of 11.0 g/cm³. This phase is preserved after cooling to room temperature, indicating an irreversible phase transformation. No formation of γ -TaON was observed during these experiments.

The calculated relative Gibbs energies as a function of the temperature (Figure 6, right, VASP/PBE results) clearly show that δ -TaON is thermodynamically unstable up to about 1600 K, and a phase transition to β -TaON should occur at lower temperatures. Experimentally, however, the δ to β phase transition is found between 1073 and 1123 K although ΔG is then much smaller compared to room temperature. Hence, kinetics is controlling the transition from δ to β .

In order to investigate the dynamic stability of δ -TaON we calculated the phonon density-of-states (PDOS) at ambient pressure. As can be seen in Figure 7 the PDOS does not exhibit any imaginary modes. Hence, the dynamic stability of the phase is assured, and anatase-type TaON is classified as a metastable solid-state material.

For the calculation of absolute band positions with respect to the vacuum, it is necessary to perform two-dimensional calculations of the TaON surface in order to circumvent the self-interaction problem in the Ewald summation. Therefore, several low-index surfaces have been investigated by performing periodic slab calculations (C09/PBE0) which are treated as truly two-dimensional systems in CRYSTAL. By calculating the surface energies E_S it was found that the (100) surface is most stable ($E_S^{\text{PBE}} = 1.1$ J/m², $E_S^{\text{PBE0}} = 1.2$ J/m²). The surface energies were converged for slab models consisting of eight or more stoichiometric layers. With the PBE0 method a

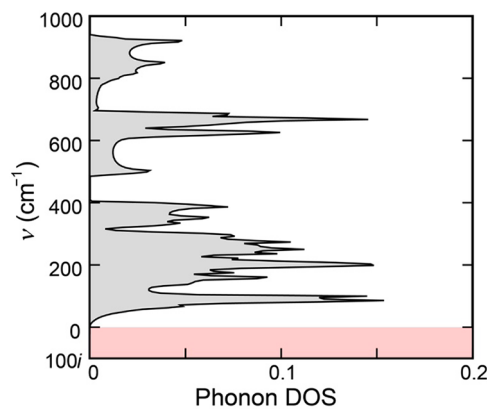


Figure 7. Phonon density of states of δ -TaON at ambient pressure calculated via density functional theory.

fundamental band gap $E_g = 3.6$ eV is obtained for the (100) surface which is slightly larger than the corresponding bulk value ($E_g = 3.2$ eV), calculated in each case as the energy difference between the highest occupied and the lowest unoccupied band. For the band gap center a value of -5.2 eV versus the vacuum level is obtained for the (100) surface of δ -TaON (C09/PBE0).

In order to obtain a value for the optical gap, we first calculated the quasiparticle gap using the GW0 approach for the δ -TaON bulk system (VASP/PBE). After five iterations where only the single particle Green's function was updated whereas the Coulomb interaction was kept fixed, we obtained a value of $E_g = 3.2$ eV for the quasiparticle gap which agrees very well with the PBE0 result for the bulk. The subsequent BSE calculation gave an excitation energy for the first optical transition of 2.7 eV which is in good agreement with the experimentally determined value for thin films of anatase-type TaON (2.4 eV).

To predict the applicability of δ -TaON as a photocatalyst in redox reactions it is necessary to determine whether tabulated redox potentials lie in between the material's valence and conduction band edges.³² Therefore, we calculated the absolute values for the valence band maximum (VBM) and the conduction band minimum (CBM) according to

$$E_{\text{VBM}} = E_{\text{BGC}}^{(100)} - \frac{1}{2}E_g^{\text{BSE}} \quad E_{\text{CBM}} = E_{\text{BGC}}^{(100)} + \frac{1}{2}E_g^{\text{BSE}} \quad (1)$$

Here E_g^{BSE} is the optical band gap in contrast to ref 32 where the quasiparticle gap was used. This way we obtained the band edge energies $E_{\text{VBM}} = -6.6$ eV and $E_{\text{CBM}} = -3.9$ eV with respect to the vacuum level. Thus, in a photocatalytic setup δ -TaON can possibly support redox reactions like the $\text{H}_2\text{O}/\text{O}_2$ couple ($E = -5.7$ eV, pH = 0) that have potentials in this range.

CONCLUSIONS

We were successful in synthesizing a new polymorph of tantalum oxide nitride with anatase-type structure by an ammonolysis reaction. The measured density is higher than the one of the thermodynamically more stable γ -TaON phase exhibiting the $\text{VO}_2(\text{B})$ structure. δ -TaON is metastable and undergoes a phase transformation to baddeleyite-type β -TaON between 800 and 850 °C. The absolute band edge positions were calculated by combining the results from periodic slab calculations with the lowest optical excitation energy for the bulk crystal, obtained with the BSE method. The resulting band edge energies ($E_{\text{VBM}} = -6.6$ eV, $E_{\text{CBM}} = -3.9$ eV) indicate that δ -TaON is a promising photocatalytic material for redox reactions, e.g., water splitting. With the enormous interest in tantalum oxide nitride phases for solar fuel production via photo- or photoelectrochemical water splitting, additional work to prepare a phase-pure material is planned. Moreover, also the synthesis of a rutile-type metastable TaON phase, which is calculated to be at most 13 kJ/mol (C09/PBE0, $T = 0$ K) less stable than the here presented anatase-type phase, is ongoing work.

ASSOCIATED CONTENT

Supporting Information

Crystallographic data in CIF format for the structure refinement of δ -TaON. Table with Gaussian-type basis sets for Ta, O, and N. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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Author Contributions

The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript.

Notes

The authors declare no competing financial interest.

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DEDICATION

Dedicated to Dr. Hans Boysen on the occasion of his 70th birthday.

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